

From the Hartree–Fock dynamics to the Vlasov equation

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switch off the external potential and study the dynamics of the low energy states

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Solution $\psi_{N,t} = e^{-iH_N t} \psi_N$, but N is huge \implies look for scaling regimes in which the evolution can be well approx. by an **effective dynamics**.

Fermionic mean-field scaling

Fix λ such that interaction and kinetic energy are of the same order:

$$E_{\text{int}} = \langle \psi_N, \lambda \sum_{i < j}^N V(x_i - x_j) \psi_N \rangle \sim \lambda N^2$$

$$E_{\text{kin}} = \langle \psi_N, \sum_{j=1}^N -\Delta_{x_j} \psi_N \rangle \sim N^{5/3}$$

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Relevant time scale: $N^{-1/3}$ (typical velocity per particle $\sim N^{1/3}$).

Rescaling time

$$i N^{1/3} \partial_t \psi_{N,t} = \left(\sum_{j=1}^N -\Delta_{x_j} + N^{-1/3} \sum_{i < j}^N V(x_i - x_j) \right) \psi_{N,t}.$$

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$$\varepsilon = N^{-1/3}$$

and multiplying by ε^2 :

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The mean-field scaling is naturally linked with a **semiclassical limit**

Effective equations

Let $\gamma_{N,t}^{(1)} = N \operatorname{tr} |\psi_{N,t}\rangle\langle\psi_{N,t}|$ be the one-particle density matrix of $\psi_{N,t}$.

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$\gamma_{N,0}^{(1)} \simeq \omega_N$ (minimiser of the Hartree-Fock energy functional)

$$\Downarrow N \gg 1$$

$\gamma_{N,t}^{(1)} \simeq \omega_{N,t}$, solution of the **Hartree-Fock equation**

$$i\varepsilon \partial_t \omega_{N,t} = [-\varepsilon^2 \Delta + V * \rho_t - X_t, \omega_{N,t}],$$

where $\rho_t(x) := \frac{1}{N} \omega_{N,t}(x; x)$ and $X_t(x; y) = \frac{1}{N} V(x - y) \omega_{N,t}(x; y)$.

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The Wigner transform of $\omega_{N,t}$

$$W_{N,t}(x, v) = \frac{\varepsilon^3}{(2\pi)^3} \int dy \omega_{N,t}\left(x + \varepsilon \frac{y}{2}; x - \varepsilon \frac{y}{2}\right) e^{-iv \cdot y}$$

for $N \rightarrow \infty$ solves the Vlasov equation

$$\partial_t W_t(x, v) + v \cdot \nabla_x W_t(x, v) = \nabla(V * \rho_t)(x) \cdot \nabla_v W_t(x, v)$$

Known results

- ★ MBQ dynamics → Vlasov:
Narnhofer and Sewell (1981), Spohn (1981).
- ★ MBQ dynamics → Hartree or Hartree–Fock:
 - ▶ Elgart, Erdős, Schlein, Yau (2004): mean–field + semiclassical scaling; analytic V , short times.
 - ▶ Benedikter, Porta and Schlein (2013): weaker assumption on V , effective estimate on the rate of convergence $\forall t$, zero temperature.
 - ▶ Petrat, Pickl (2014): similar result with a different method.
 - ▶ Benedikter, Porta, Jakšić, S. and Schlein (2015): positive temperature.
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 - ▶ Golse, Paul (2015): ∇V Lipschitz continuous, explicit estimates on the rate of convergence.

Notations

ω_N sequence of reduced densities on $L^2(\mathbb{R}^3)$

$$\text{tr } \omega_N = N, \quad 0 \leq \omega_N \leq 1$$



Wigner trans.

W_N

Hartree eq.n



$\omega_{N,t}$

$$i\varepsilon \partial_t \omega_{N,t} = [h(t), \omega_{N,t}]$$

$$h(t) = -\varepsilon^2 \Delta + V * \rho_t$$

Vlasov eq.n

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$\tilde{W}_{N,t}$

Weyl quantization



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$$A_t(x; y) = \nabla(V * \tilde{\rho}_t) \left(\frac{x+y}{2} \right) \cdot (x-y) \tilde{\omega}_{N,t}(x; y)$$

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- $\|W_N\|_{H_a^s}^2 = \sum_{\beta \leq s} \int (1 + x^2 + v^2)^a |\nabla^\beta W_N(x, v)|^2 dx dv$

Statement of the results: mixed states

Theorem 1 (Benedikter, Porta, S., Schlein)

Let $V \in \mathcal{W}^{2,\infty}(\mathbb{R}^3)$ and $\|W_N\|_{H_4^5} < C$.

Then there exists $C = C(\|V\|_{\mathcal{W}^{2,\infty}}, \sup_N \|W_N\|_{H_4^2}) > 0$ such that

$$\mathrm{tr} |\omega_{N,t} - \tilde{\omega}_{N,t}| \leq CN\varepsilon \exp(C \exp(C|t|)) [1 + \sum_{k=1}^3 \varepsilon^k \sup_N \|W_N\|_{H_4^{k+2}}].$$

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Theorem 2 (Benedikter, Porta, S., Schlein)

Let $V \in L^1(\mathbb{R}^3)$ such that $A = \int |\hat{V}(p)| (1 + |p|^2) dp < \infty$ and $\|W_N\|_{H_4^2} < C$.

Then there exists $C = C(\|W_N\|_{H_4^2}, A) > 0$ such that

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Equivalently,

$$\|W_{N,t} - \tilde{W}_{N,t}\|_{L^2} \leq C\varepsilon \exp(C \exp(C|t|)).$$

Statement of the results: pure states

At zero temperature we prove convergence for the expectation of a class of **semiclassical observables**, functions of x and $-i\varepsilon\nabla$.

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Let $V \in L^1(\mathbb{R}^3)$ be such that $\int |\hat{V}(p)|(1 + |p|^3) dp < \infty$.

Let ω_N satisfy

$$\mathrm{tr} |[x, \omega_N]| \leq CN\varepsilon, \quad \mathrm{tr} |[\varepsilon\nabla, \omega_N]| \leq CN\varepsilon.$$

Let $W_N \in L^1(\mathbb{R}^3 \times \mathbb{R}^3)$ be such that $\|W_N\|_{\mathcal{W}^{1,1}} \leq C$.

Then there exists $C > 0$ such that

$$|\mathrm{tr} e^{ip \cdot x + \varepsilon q \cdot \nabla} (\omega_{N,t} - \tilde{\omega}_{N,t})| \leq CN\varepsilon (1 + |p| + |q|)^2 \exp(C|t|),$$

$$\forall p, q \in \mathbb{R}^3, \quad \forall t \in \mathbb{R}.$$

Strategy: trace norm convergence

Compare $\omega_{N,t}$ and $\tilde{\omega}_{N,t}$ via Gronwall inequality

Step 1.

$$i\varepsilon \partial_t (\omega_{N,t} - \tilde{\omega}_{N,t}) = [-\varepsilon^2 \Delta, (\omega_{N,t} - \tilde{\omega}_{N,t})] + \dots$$

how to get rid of the kinetic term?

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Define a unitary operator \mathcal{U} ,

$$\begin{cases} i\varepsilon \partial_t \mathcal{U}(t) = h(t) \mathcal{U}(t), \\ \mathcal{U}(0) = 1. \end{cases}$$

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Then

$$\begin{aligned} i\varepsilon\partial_t\mathcal{U}^*(t)(\omega_{N,t} - \tilde{\omega}_{N,t})\mathcal{U}(t) \\ &= -\mathcal{U}^*(t)[h(t), \omega_{N,t} - \tilde{\omega}_{N,t}]\mathcal{U}(t) \\ &\quad + \mathcal{U}^*(t)([h(t), \omega_{N,t} - \tilde{\omega}_{N,t}] + [V * \rho_t, \tilde{\omega}_{N,t}] - A_t)\mathcal{U}(t) \\ &= \mathcal{U}^*(t)([V * (\rho_t - \tilde{\rho}_t), \tilde{\omega}_{N,t}] + [V * \tilde{\rho}_t, \tilde{\omega}_{N,t}] - A_t)\mathcal{U}(t). \end{aligned}$$

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where

$$B_t(x; y) = [V * \tilde{\rho}_t(x) - V * \tilde{\rho}_t(y) - \nabla(V * \tilde{\rho}_t)\left(\frac{x+y}{2}\right) \cdot (x-y)] \tilde{\omega}_{N,t}(x; y)$$

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Step 2.

Integrate in time and take the trace norm:

$$\begin{aligned} \|\omega_{N,t} - \tilde{\omega}_{N,t}\|_{\text{Tr}} &\leq \|\omega_{N,0} - \tilde{\omega}_{N,0}\|_{\text{Tr}} \\ &\quad + \frac{1}{\varepsilon} \int_0^t \| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} ds \\ &\quad + \frac{1}{\varepsilon} \int_0^t \|B_s\|_{\text{Tr}} ds \end{aligned}$$

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$$\frac{1}{\varepsilon} \int_0^t \| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} ds$$

$$\| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} \leq \| \rho_s - \tilde{\rho}_s \|_{L^1} \sup_z \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}}$$

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$$\frac{1}{\varepsilon} \int_0^t \| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} ds$$

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heuristically

$$\begin{aligned} \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}} & \text{ "}\leq\text{"} \| [x, \tilde{\omega}_{N,s}] \|_{\text{Tr}} \text{ "}= " N\varepsilon \| \nabla \tilde{W}_{N,s} \|_{L^1_{x,v}} \\ & \leq C N\varepsilon \| \nabla W_{N,0} \|_{L^1} e^{C|s|} \end{aligned}$$

Strategy

$$\frac{1}{\varepsilon} \int_0^t \| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} ds$$

$$\| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} \leq \| \rho_s - \tilde{\rho}_s \|_{L^1} \sup_z \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}}$$

heuristically

$$\begin{aligned} \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}} & \leq \| [x, \tilde{\omega}_{N,s}] \|_{\text{Tr}} = N\varepsilon \| \nabla \tilde{W}_{N,s} \|_{L^1_{x,v}} \\ & \leq C N\varepsilon \| \nabla W_{N,0} \|_{L^1} e^{C|s|} \end{aligned}$$

$$\| \rho_s - \tilde{\rho}_s \|_{L^1} \leq \frac{1}{N} \| \omega_{N,s} - \tilde{\omega}_{N,s} \|_{\text{Tr}}$$

Strategy

$$\frac{1}{\varepsilon} \int_0^t \| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} ds$$

$$\| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} \leq \| \rho_s - \tilde{\rho}_s \|_{L^1} \sup_z \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}}$$

heuristically

$$\begin{aligned} \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}} & \leq \| [x, \tilde{\omega}_{N,s}] \|_{\text{Tr}} = N\varepsilon \| \nabla \tilde{W}_{N,s} \|_{L^1_{x,v}} \\ & \leq C N\varepsilon \| \nabla W_{N,0} \|_{L^1} e^{C|s|} \end{aligned}$$

$$\| \rho_s - \tilde{\rho}_s \|_{L^1} \leq \frac{1}{N} \| \omega_{N,s} - \tilde{\omega}_{N,s} \|_{\text{Tr}}$$

$$\begin{aligned} \| [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{Tr}} & \leq \| (1 - \varepsilon^2 \Delta)^{-1} (1 + x^2)^{-1} \|_{\text{HS}} \| (1 + x^2) (1 - \varepsilon^2 \Delta) [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{HS}} \\ & \leq C \sqrt{N} \| (1 + x^2) (1 - \varepsilon^2 \Delta) [V(z - \cdot), \tilde{\omega}_{N,s}] \|_{\text{HS}} \\ & \leq C \exp(C|s|) N\varepsilon [\| W_N \|_{H^1_4} + \varepsilon \| W_N \|_{H^2_4} + \varepsilon^2 \| W_N \|_{H^3_4} + \varepsilon^3 \| W_N \|_{H^4_4}] \end{aligned}$$

Strategy

$$\begin{aligned}\|\omega_{N,t} - \tilde{\omega}_{N,t}\|_{\text{Tr}} &\leq \|\omega_{N,0} - \tilde{\omega}_{N,0}\|_{\text{Tr}} \\ &+ \frac{1}{\varepsilon} \int_0^t \| [V * (\rho_s - \tilde{\rho}_s), \tilde{\omega}_{N,s}] \|_{\text{Tr}} ds \\ &+ \frac{1}{\varepsilon} \int_0^t \| B_s \|_{\text{Tr}} ds\end{aligned}$$

Strategy

$$\frac{1}{\varepsilon} \int_0^t \|B_s\|_{\text{Tr}} ds$$

with

$$B_t(x; y) = [V * \tilde{\rho}_t(x) - V * \tilde{\rho}_t(y) - \nabla(V * \tilde{\rho}_t)\left(\frac{x+y}{2}\right) \cdot (x-y)] \tilde{\omega}_{N,t}(x; y)$$

Strategy

$$\frac{1}{\varepsilon} \int_0^t \|B_s\|_{\text{Tr}} ds$$

with

$$B_t(x; y) = [V * \tilde{\rho}_t(x) - V * \tilde{\rho}_t(y) - \nabla(V * \tilde{\rho}_t)\left(\frac{x+y}{2}\right) \cdot (x-y)] \tilde{\omega}_{N,t}(x; y)$$

$$\|[x, [x, \tilde{\omega}_{N,s}]]\|_{\text{Tr}} = N \varepsilon^2 \|\nabla^2 \tilde{W}_s\|_{L^1_{x,v}}$$

Strategy

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$$\|[x, [x, \tilde{\omega}_{N,s}]]\|_{\text{Tr}} = N \varepsilon^2 \|\nabla^2 \tilde{W}_s\|_{L^1_{x,v}}$$

Strategy

$$\frac{1}{\varepsilon} \int_0^t \|B_s \tilde{\omega}_{N,s}\|_{\text{Tr}} ds$$

with

$$B_t(x; y) = V * \tilde{\rho}_t(x) - V * \tilde{\rho}_t(y) - \nabla(V * \tilde{\rho}_t)\left(\frac{x+y}{2}\right)$$

$$\|[x, [x, \tilde{\omega}_{N,s}]]\|_{\text{Tr}} = N \varepsilon^2 \|\nabla^2 \tilde{W}_s\|_{L^1_{x,v}}$$

Hence we conclude by Gronwall Lemma:

$$\begin{aligned} \|\omega_{N,t} - \tilde{\omega}_{N,t}\|_{\text{Tr}} &\leq C \int_0^t ds \|\omega_{N,s} - \tilde{\omega}_{N,s}\|_{\text{Tr}} \exp(C|s|) \\ &+ C N \varepsilon [\|W_N\|_{H_4^2} + \varepsilon \|W_N\|_{H_4^3} + \varepsilon^2 \|W_N\|_{H_4^4} + \varepsilon^3 \|W_N\|_{H_4^5}] \exp(C|s|) \end{aligned}$$